

The deviation of the point for specimen 581 from the straight line which fits the rest of the data is not understood at present, but it seems fairly certain that this deviation does not represent experimental error. It is of interest to note that specimen 581 was also exceptional in exhibiting a much broader transition than any of the other specimens. This breadth of transition as well as the unusually large critical field shift shown in Fig. 1 have been observed with excellent reproducibility in two runs between which the specimen was repurified and recrystallized. The only respect in which specimen 581 appears to differ markedly from the others is that it has a relatively broad mass distribution whereas the remaining specimens have their dominant isotopic percentages situated fairly close to M . This anomalous behavior will receive further attention.

Measurements are now being extended down to about 1.2°K to investigate the temperature dependence of the critical field shift. A more detailed account of this work will be prepared when this phase of the measurements has been completed.

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¹ H. Fröhlich, *Phys. Rev.* **79**, 845 (1950).

² J. Bardeen, *Phys. Rev.* **80**, 567 (1950).

³ E. Maxwell, *Physics Today* **5**, 14 (1952), or more recently, B. Serin, *Progress in Low Temperature Physics*, edited by C. J. Gorter (Interscience Publishers, Inc., New York, 1955), Vol. I, p. 142.

⁴ M. Olsen, *Nature* **168**, 245 (1951).

⁵ Serin, Reynolds, and Lohman, *Phys. Rev.* **86**, 162 (1951).

⁶ For a theoretical discussion of the lead result see J. de Launay, *Phys. Rev.* **93**, 661 (1954).

⁷ Cochran, Mapother, and Mould, *Phys. Rev.* **103**, 1657 (1956).

⁸ R. R. Hake and D. E. Mapother, *Intern. J. Phys. Chem. Solids* (to be published).

⁹ Boorse, Cook, and Zemansky, *Phys. Rev.* **78**, 635 (1950).

¹⁰ $(\partial H_c/\partial T)_M$ is the slope of the critical field curve at T_c and is computed from the specific heat data of J. R. Clement and E. H. Quinell, *Phys. Rev.* **85**, 502 (1952). It is in agreement with the critical field data of Daunt, Horseman, and Mendelssohn, *Phil. Mag.* **27**, 754 (1939).

Metallic Transition in Ionic and Molecular Crystals

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EXPERIMENTAL evidence has been obtained that several ionic and molecular crystals under pressures less than 250 000 atmospheres undergo a transition to a metallic state as indicated by their conductivity. The pressure is created by a shock wave

generated from a high-explosive system. A decrease in the resistance by a factor of at least 10^6 from the uncompressed to the compressed state is measured by means of pins¹ pressed against the surface of the crystal.

The materials that showed a resistance of less than a few hundred ohms at 250 000 atmospheres were a single crystal of CsI and the following powders which were compressed into pellets to densities near that of the crystal; commercial grade I_2 , red phosphorus, and $LiAlH_4$.² Teflon and pressed NaCl pellets did not show a substantial decrease in resistance. It is believed that this confirms the validity of the experimental procedure. Another substance, compressed CsBr powder, appeared to have a resistance of the order of kilo-ohms, which might mean that it is in the transition region from an insulator to a conductor. Initially all samples had a resistance greater than 10^8 ohms except red phosphorus, whose resistance was about 5×10^6 ohms.

The high-explosive system is similar to that previously used.³ This assembly induces a strong plane shock into a series of plates in contact with the high explosive. This series of plates consists of three aluminum plates, individually grounded, and separated by two Teflon plates. This is believed to produce an uncharged shock at the sample. Each sample is pressed against the last plate by two spring-loaded pins. One of these pins is charged to 300 volts while the other is grounded. These pins constitute part of an appropriate RC circuit. A "raster"-type oscilloscope is attached across a resistor of this circuit, so that any discharge can be observed and photographed. The shape of the signal indicates the resistance of the sample. More precise resistance measurements are being worked out.

The pressure is determined from the known equation of state of aluminum⁴ by means of an accurate measurement of the velocity of the surface. This is accomplished by a group of seven pins, spaced at 1.5-mm intervals away from the aluminum plate.

An apparent shock velocity is obtained by measuring the time difference between the signal from the pins pressed into the face of the sample and the closure of an additional pin located beside the base of the sample and 0.2 mm away from the surface of the aluminum

TABLE I. Shock velocity (mm/ μ sec).

Sample	Expt. 1	Expt. 2	Known
CsBr	4.3	...	4.1 ± 0.2^b
CsI	...	3.8	3.8 ± 0.2^b
I_2	3.7	3.7	°
$LiAlH_4$	5.6	5.5	°
P(red)	3.8	4.0	5.0 ± 1.5^d
NaCl	2.0^a	2.3^a	5.9 ± 0.2^b
Teflon	...	1.6^a	°

^a These low values indicate that the aluminum plate rather than a shock-induced metallic transition discharged these pins.

^b Calculated from equation of state experiments performed at Livermore.

^c Any velocity greater than 2.8 mm/ μ sec, the measured velocity of the aluminum surface, cannot be the result of the arrival of this surface at the pins.

^d Estimated from compressibility data, P. W. Bridgman, *The Physics of High Pressures* (G. Bell and Sons, London, 1949).